

Bubble Growth in Cosmological 1st Order Phase Transitions: Runaway Walls vs Hydrodynamic Obstructions

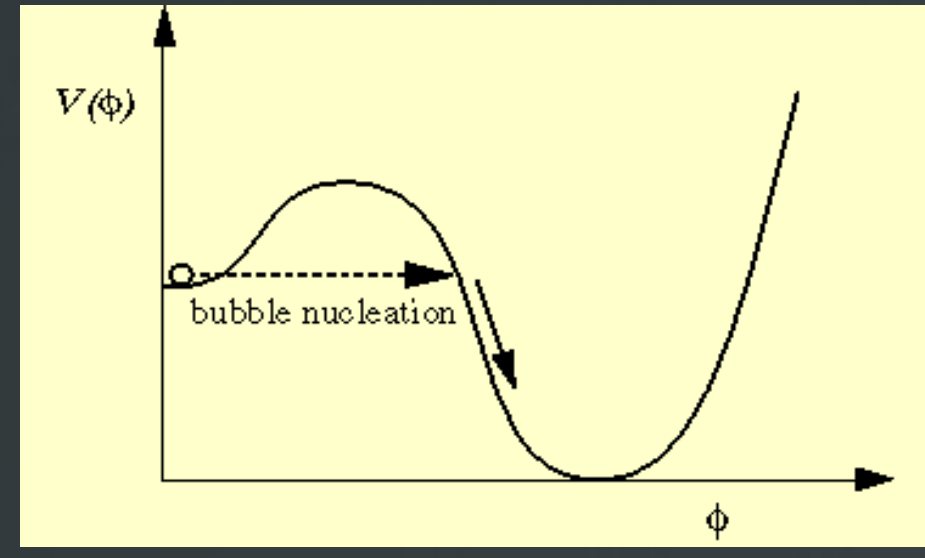
J. R. Espinosa, T. Konstandin, J. M. No and G. Servant, JCAP **1006** (2010) 028

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Motivation

Why Bubbles? → Phase Transitions in the Early Universe

If 1st Order...



Nucleation & Growth of True Vacuum Bubbles in False Vacuum Sea.
(Nucleation Temperature T_N)



Bubble Growth in Cosmological Phase Transitions Relevant for:

- **Electroweak Baryogenesis**
1st Order Phase Transition = 3rd Sakharov condition.
Viable Baryogenesis → $v_w < c_s$ ($v_w \ll 1$ is Favoured)
- **Stochastic Background of Gravitational Waves**
Bubble Collisions ($\propto v_w^4$) + Turbulence in Plasma

Formalism

① Matching Equations Across the Bubble Wall.

System “Higgs Wall – Plasma” ⇒
$$\begin{aligned} T_{\mu\nu}^\phi &= \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left(\frac{1}{2} \partial_\rho \phi \partial^\rho \phi - V_0(\phi) \right) \\ T_{\mu\nu}^{\text{Plasma}} &= w u_\mu u_\nu - g_{\mu\nu} p \end{aligned}$$

Energy-Momentum Conserved Across Wall
+
Steady State (Wall Reference Frame)

$$\begin{aligned} w_+ v_+^2 \gamma_+^2 + p_+ - V_0(0) &= w_- v_-^2 \gamma_-^2 + p_- - V_0(\phi_0) \\ w_+ v_+ \gamma_+^2 &= w_- v_- \gamma_-^2 \end{aligned}$$

With $V_0(0) - V_0(\phi_0) \equiv \epsilon$

$$v_+ = \frac{1}{1 + \alpha_+} \left(\frac{v_-}{2} + \frac{1}{6v_-} \pm \sqrt{\left(\frac{v_-}{2} + \frac{1}{6v_-} \right)^2 + \alpha_+^2 + \frac{2}{3}\alpha_+ - \frac{1}{3}} \right)$$

$$\alpha_+ \equiv \frac{\epsilon}{a_+ T_+^4}$$

Wall

$\phi = 0$ $\phi = \phi_0$

v_+ v_-

T_+ T_-

④ Higgs Equation of Motion. (Needed to fix ξ_w)

$$\partial_\mu \partial^\mu \phi + \frac{\partial \mathcal{F}}{\partial \phi} - \mathcal{K}(\phi) = 0$$

Phenomenological Ansatz

$$\mathcal{K}(\phi) = T_N \tilde{\eta} \frac{u^\mu \partial_\mu \phi}{\sqrt{1 + (\lambda_\mu u^\mu)^2}}$$

With $\mathcal{K}(\phi) = - \sum_i \frac{dm_i^2}{d\phi} \int \frac{d^3 p}{(2\pi)^3 2E_i} \delta f_i(p)$

$$\partial_z^2 \phi - \frac{\partial \mathcal{F}}{\partial \phi} = T_N \tilde{\eta} v \partial_z \phi \rightarrow \alpha_+ - \alpha_c = \eta \frac{\alpha_+}{\alpha_N} \langle v \rangle$$

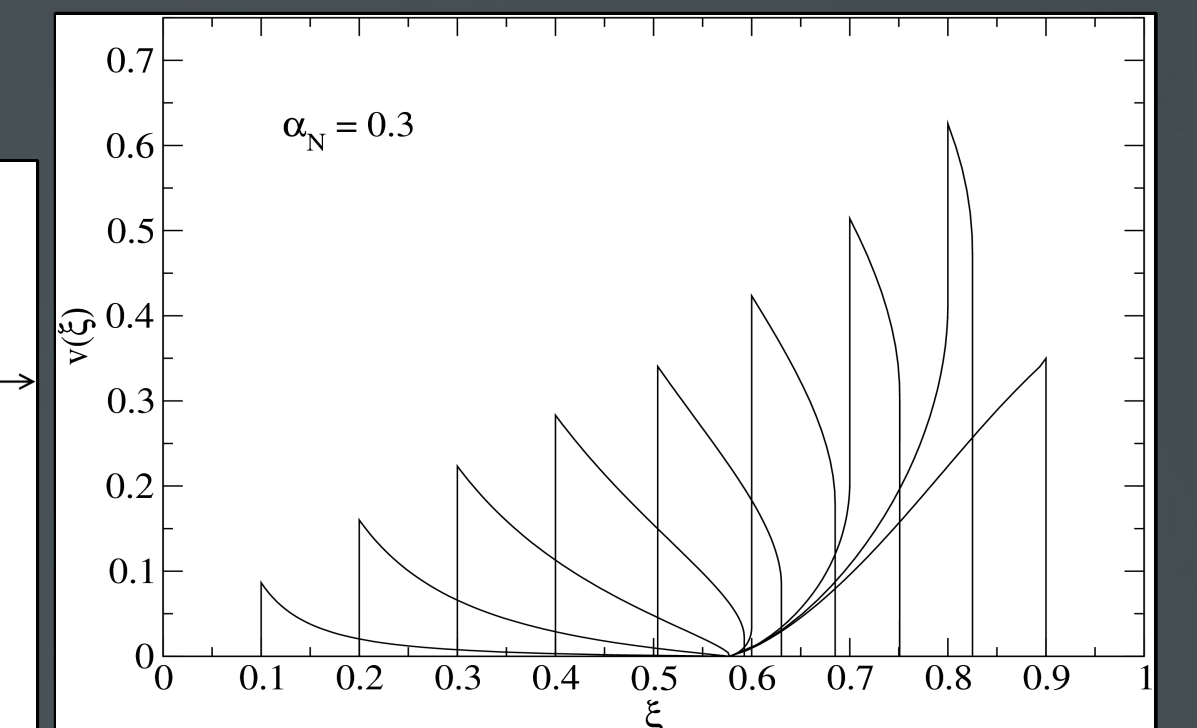
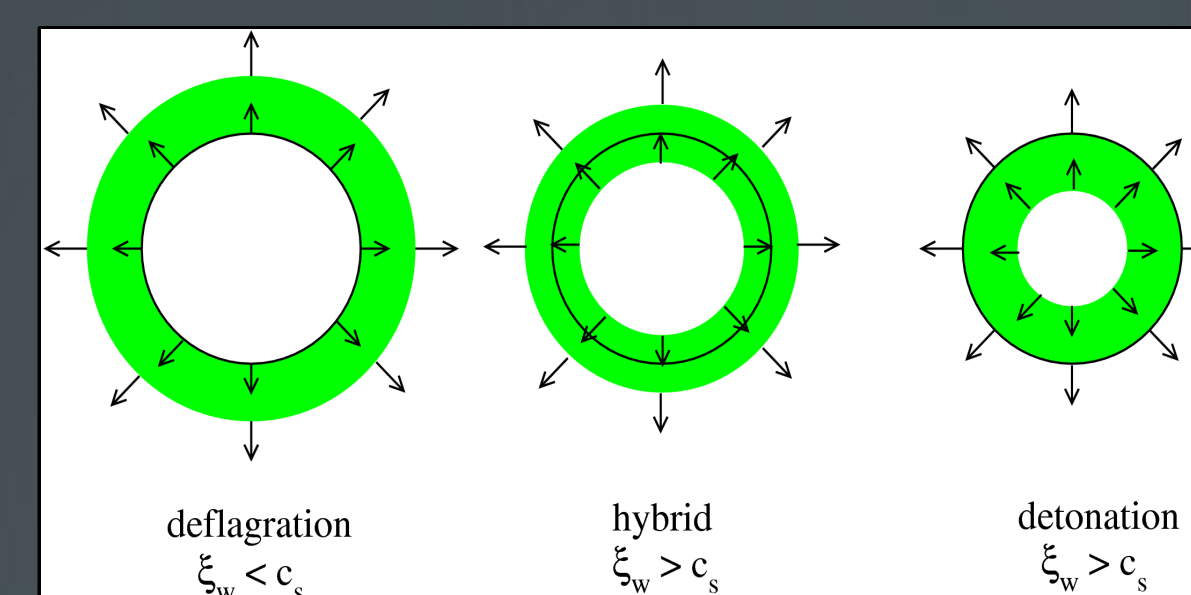
η Friction Parameter

② Fluid Equations for Plasma. $\left(\frac{\partial_\mu T_{\text{Plasma}}^{\mu\nu} = 0}{v(r, t) = v(\xi = r/t)} \right)$

$$\frac{1 - \xi v(\xi)}{1 - v^2(\xi)} \left[\frac{\mu^2}{c_s^2} - 1 \right] \partial_\xi v = 2 \frac{v(\xi)}{\xi} + w(\xi) = w_0 \exp \left[\left(1 + \frac{1}{c_s^2} \right) \int_{v_0}^{v(\xi)} \gamma^2 \mu dv \right]$$

③ Fluid Solutions.

- **Deflagrations:** Subsonic v_w . Fluid at Rest Behind Bubble Wall.
 $c_s > v_- = v_w > v_+$ Compression Wave in Front of Wall. $T_+ > T_-$
- **Detonations:** Supersonic v_w . Fluid at Rest in Front of Wall.
 $v_w = v_+ > v_- > c_s$ Rarefaction Wave Behind Wall. $T_+ = T_-$
- **Hybrids.**



Hydrodynamic Obstruction

Higgs EoM:

Steady State Solution

$$\partial_\mu \partial^\mu \phi + \frac{\partial \mathcal{F}}{\partial \phi} - \mathcal{K}(\phi) = 0 \rightarrow F_{dr} = \int dz \partial_z \phi \frac{\partial \mathcal{F}}{\partial \phi} = \int dz \partial_z \phi \mathcal{K}(\phi) = F_{fr}$$

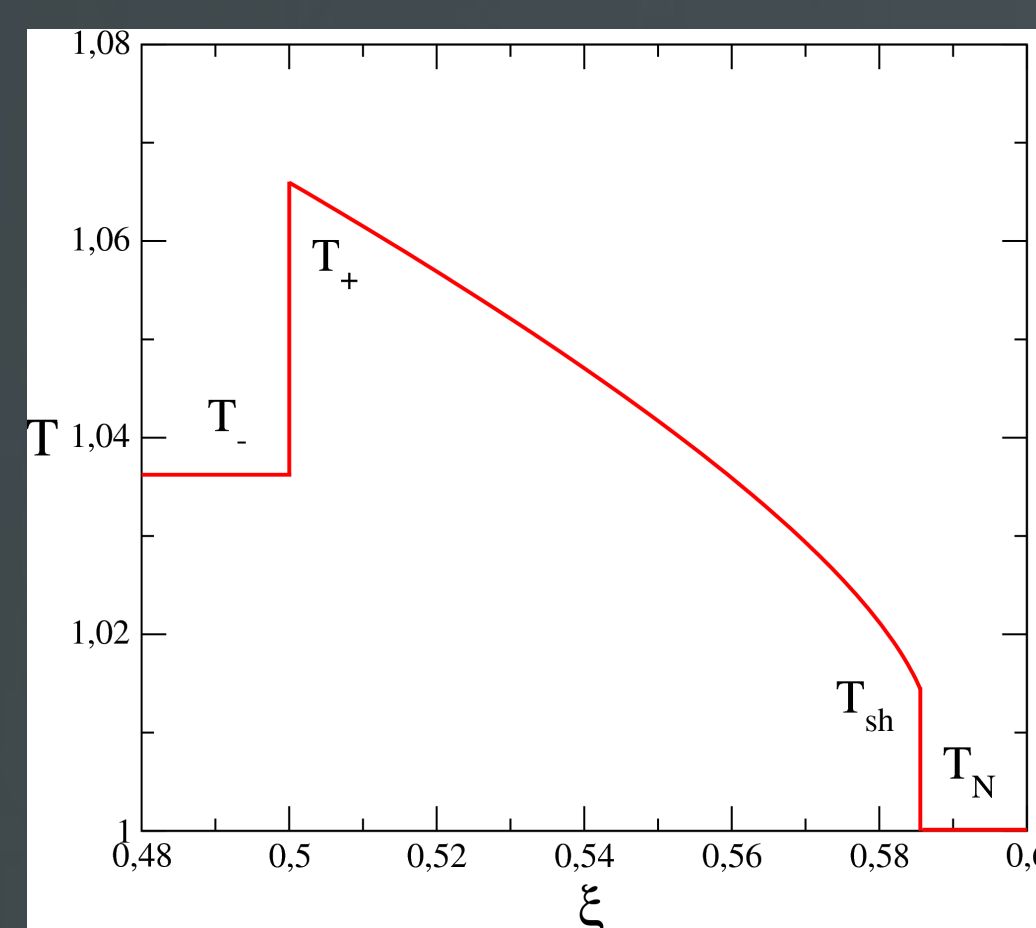
Suppose Deflagration Solution:

- ⇒ As Wall Moves, Reheats Plasma in Front ($T_+ > T_-$).
- ⇒ As v_w Increases, T_+ Raises ($T_+ \uparrow$ if $v_w \uparrow$).
- ⇒ As T_+ Raises, F_d Decreases ($F_d \downarrow$ if $T_+ \uparrow$).



It could happen that $F_d \rightarrow 0$ for $v_w < c_s$

Hydrodynamic Obstruction!!



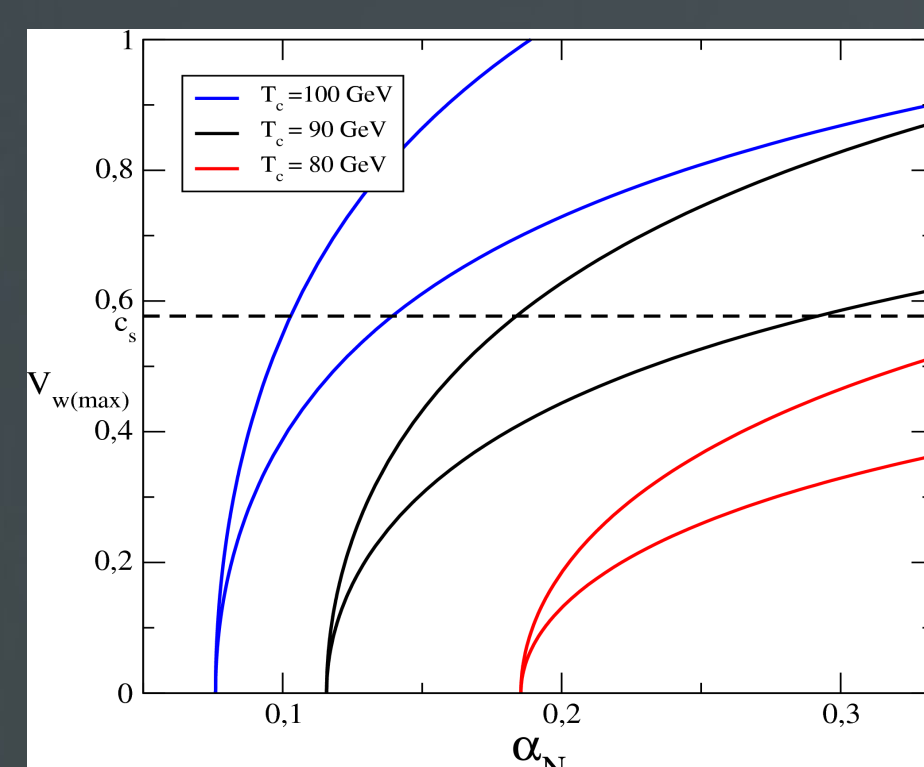
Criterion for Hydrodynamic Obstruction:

$$\alpha_N < \alpha_c \text{Exp} \left(\frac{16}{3(1 + \alpha_c)} \Omega(\alpha_c) \right) \quad \text{With} \quad \Omega(\alpha_c) \equiv \sqrt{\frac{\alpha_c}{2} - \frac{3}{10}\alpha_c - \frac{1}{5}\alpha_c^3/2} \quad v_w = c_s$$

Saturates Inequality

If Obstruction Occurs, $T_+ > T_c > T_-$

- **Lower Limit** ($T_+ = T_c$): $v_w^{\text{max}} = \sqrt{\frac{1}{6\alpha_c} \text{Log} \left(\frac{T_c}{T_N} \right)}$
($v_w \ll 1$ approximation)
- **Upper Limit** ($T_- = T_c$): $v_w^{\text{max}} = \sqrt{\frac{1}{3\alpha_c} \text{Log} \left(\frac{T_c}{T_N} \right)}$



Higgs EoM Revisited:

$$\dot{v}_w \propto (\mathcal{F}_- - \mathcal{F}_+)(T) - v_w \eta \rightarrow \dot{v}_w \propto (\mathcal{F}_- - \mathcal{F}_+)(T_N) - v_w^2 \kappa - v_w \eta$$

Even for Vanishing Friction, Small v_w Still Possible (depends on α_N).
(Hydrodynamic Obstruction to Larger v_w)

Runaway Solutions Change Energy Budget of 1st Order Phase Transitions.
(Also, Detonation Parameter Space Drastically Reduced).

Runaway Walls

Higgs EoM:

$$\partial_\mu \partial^\mu \phi + \frac{\partial \mathcal{F}}{\partial \phi} - \mathcal{K}(\phi) = 0 \quad F_{fr} = \int dz \partial_z \phi \mathcal{K}(\phi) \quad F_{dr} = \int dz \partial_z \phi \frac{\partial \mathcal{F}}{\partial \phi}$$

Driving Force Friction Force

⇒ For $v_w \ll 1$ (Low Velocity Limit) → $F_{fr} \propto \gamma_w v_w \sim v_w$

⇒ For $v_w \rightarrow 1$ (Ultrarelativistic Limit) → $F_{fr} \rightarrow F_{fr}^{\text{max}} = Cte$

D. Bodecker and G. D. Moore, JCAP **0905** (2009) 009

• If $F_d < F_{fr}^{\text{max}}$:

Friction Equilibrates with Driving Force.

Steady State Solution.

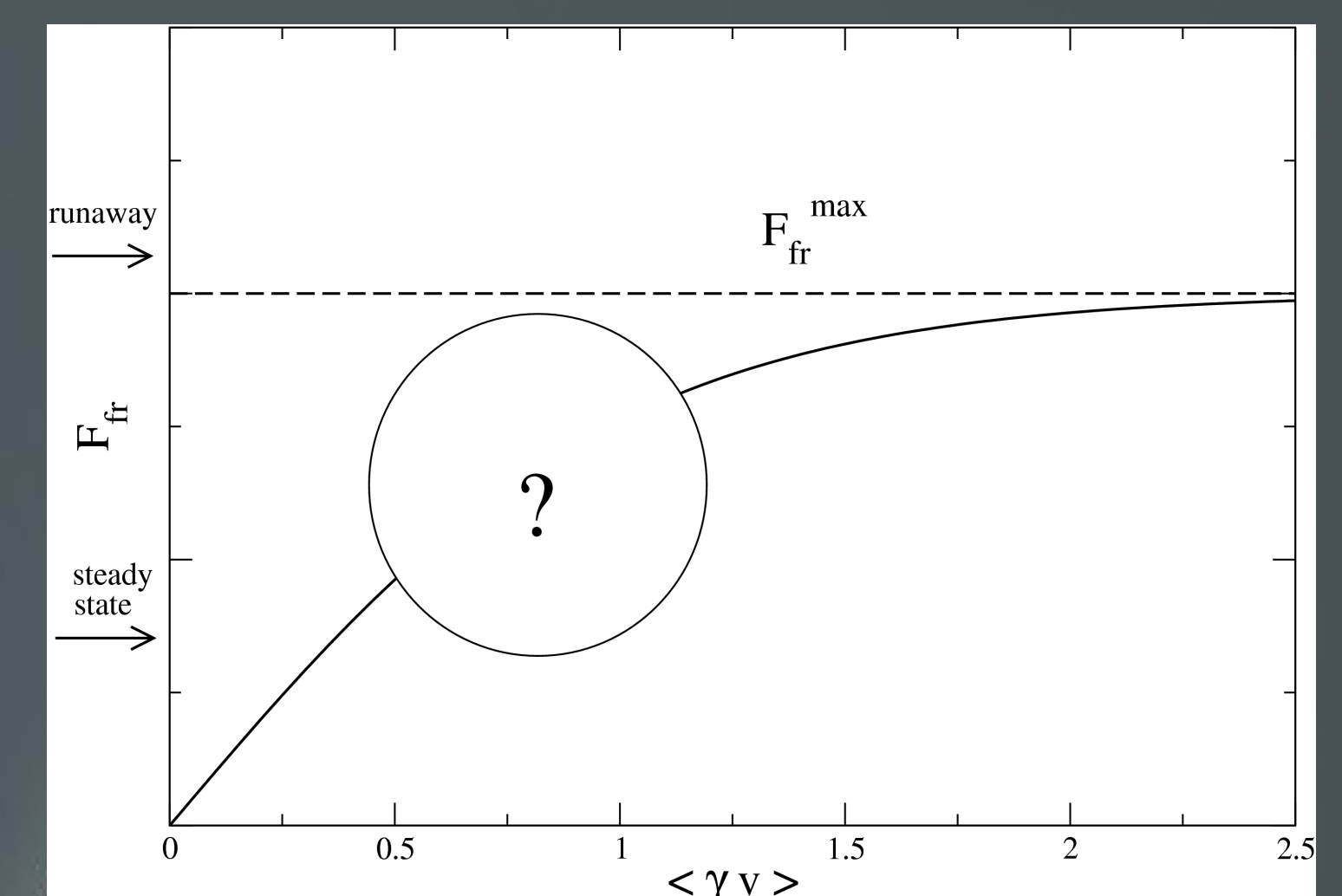
(Deflagrations / Hybrids / Detonations)

• If $F_d > F_{fr}^{\text{max}}$:

Friction does not Equilibrate Driving Force.

Runaway Solution.

Continuous Acceleration (No Steady State)

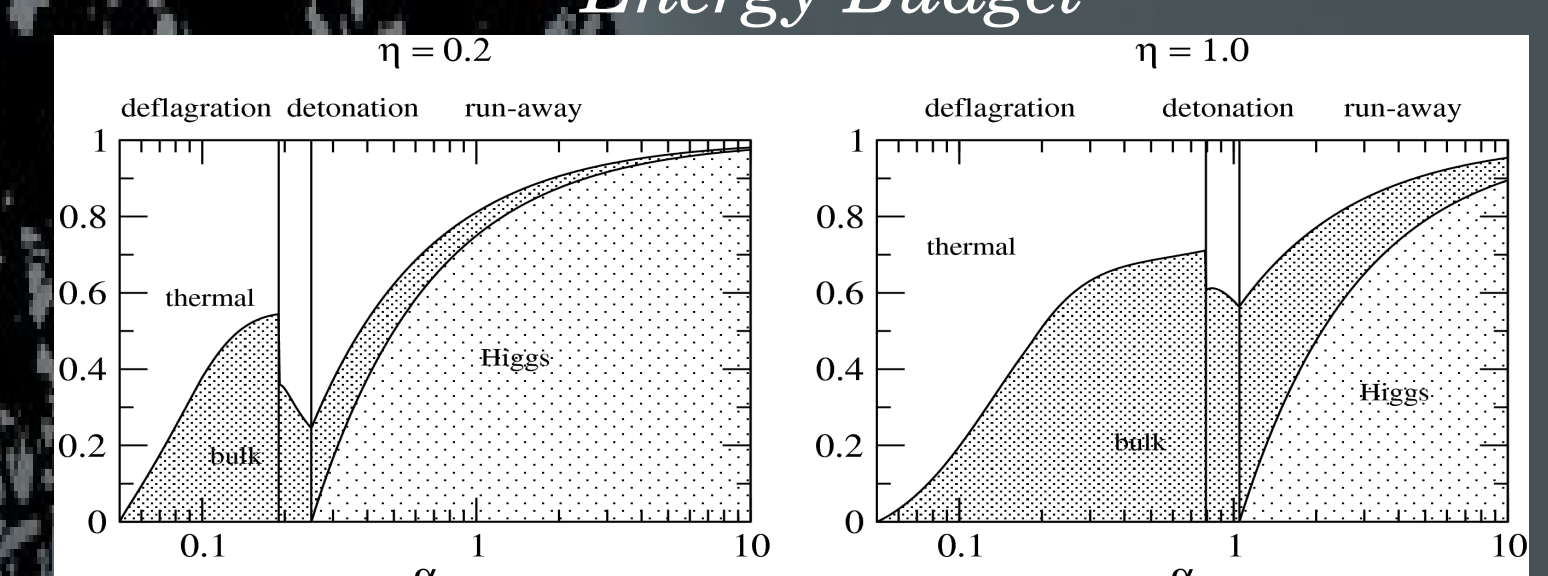


Runaway Criterion ($F_d > F_{fr}^{\text{max}}$)

$$\alpha_N > \alpha_\infty \equiv \frac{30}{\pi^2} \left(\frac{\phi_0}{T_N} \right)^2 \frac{\sum_{\text{light} \rightarrow \text{heavy}} c_i \|N_i\| y_i^2}{\sum_{\text{light}} c_i' \|N_i\|}$$

If $\alpha_N > \alpha_\infty$, Some Energy into Higgs Field
Energy Budget Altered!!

Energy Budget



Conclusions

⇒ Favours Baryogenesis

⇒ May Change Gravitational Wave Spectrum (Mainly Turbulence)